

Test 3 Solutions

1) Before transfer: $E = -\frac{k}{2a} = -\frac{k}{2r_1} = \frac{\mu v^2}{2} - \frac{k}{r_1} \Rightarrow \frac{\mu v^2}{2} = \frac{k}{2r_1}$
 $\Rightarrow v = \sqrt{k/\mu r_1}$

After transfer $E = -\frac{k}{2a} = -\frac{k}{(r_1+r_2)} = \frac{\mu v_{f1}^2}{2} - \frac{k}{r_1} \Rightarrow v_{f1} = \sqrt{\frac{2kr_2}{\mu(r_1+r_2)}}$

2) Consider a rotating coordinate system at the center of the earth whose axes are parallel to the local axes as follows: x : East, y : North, z : up.

Then $\vec{\omega} = \omega(\cos\lambda \hat{y} + \sin\lambda \hat{z})$, $\vec{r} \approx R \hat{z}$, $\vec{v} \approx -gt \hat{z}$

Without corrections the motion is described by

$z = (h - gt^2)/2$ and $x = y = 0$. The small accelerations are

$\vec{a}_{\text{cen.}} = -\vec{\omega} \times (\vec{\omega} \times \vec{r}) = -\omega^2 \hat{\omega} \times (R \cos\lambda \hat{x}) = \omega^2 R (\cos\lambda \hat{z} - \sin\lambda \hat{y})$, and

$\vec{a}_{\text{cor.}} = -2\vec{\omega} \times \vec{v} \approx +2\omega \cos\lambda gt \hat{x}$

Neglecting the z -component of $\vec{a}_{\text{cen.}}$ we get for the deflection two significant contributions:

Centrifugal $\vec{a}_{\text{cen.}} \approx -\omega^2 R \sin\lambda \frac{T^2}{2} \hat{y} = \frac{\omega^2 R h \sin\lambda}{g}$ (towards south)

Coriolis $\vec{a}_{\text{cor.}} \approx 2\omega \cos\lambda g \frac{T^3}{6} \hat{x} = \frac{\omega \cos\lambda}{3} \left(\frac{8h^3}{g}\right)^{1/2}$ (towards east)

Which of these is larger in magnitude depends on h and λ : the Coriolis effect dominates if $\cos\lambda \sqrt{h} \gg \frac{\omega R \sin\lambda}{\sqrt{g}}$
 i.e., if $\sqrt{gh} \gg \omega R \tan\lambda$.

$$\#3)(a) \quad I_{ij} = \sum_{\alpha} m_{\alpha} (r_{\alpha}^2 \delta_{ij} - r_{\alpha,i} r_{\alpha,j})$$

For the O-atom all coordinate components are zero.

Thus, we need only include the H-atoms:

$$I_{ij} = m \left\{ 1 \delta_{ij} - (1,0,0)(1,0,0) + (1 \delta_{ij} - (-0.25, 0.97, 0)(-0.25, 0.97, 0)) \right\}$$

$$= m \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \overset{\circ}{\text{Å}}^2 + m \begin{bmatrix} 15/16 & 0.24 & 0 \\ 0.24 & 1/16 & 0 \\ 0 & 0 & 1 \end{bmatrix} \overset{\circ}{\text{Å}}^2$$

$$= m \begin{bmatrix} 0.94 & 0.24 & 0. \\ 0.24 & 1.06 & 0. \\ 0. & 0. & 2 \end{bmatrix} \overset{\circ}{\text{Å}}^2$$

$$(b) \quad \vec{R}_{CM} = \sum_{\alpha} \frac{m_{\alpha} \vec{r}_{\alpha}}{M} = \frac{(1,0,0)}{18} + \frac{(-0.25, 0.97, 0.)}{18} = \frac{(0.75, 0.97, 0.)}{18}$$

$$= (0.04, 0.055, 0.) \overset{\circ}{\text{Å}}$$

$$\text{Thus, } J_{ij} = I_{ij} - 18m(R^2 \delta_{ij} - R_i R_j)$$

Notice that the O-atom is so heavy that up to the second decimal place in $\overset{\circ}{\text{Å}}^2$ the correction terms $\sim R^2$ are all zero, and $J_{ij} \approx I_{ij}$.